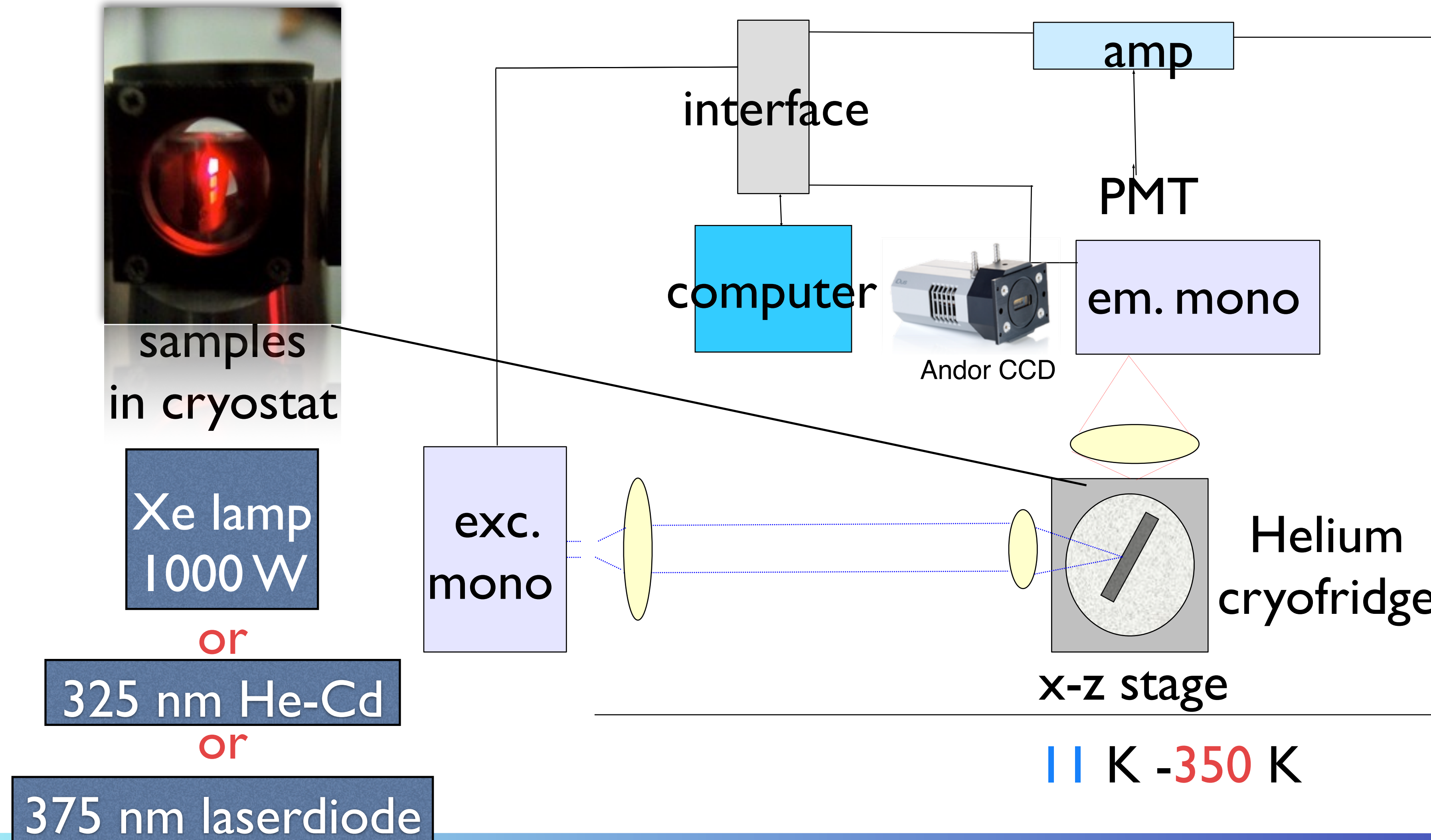


IQE and all that jazz: the temperature dependence of semiconductor light emission

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For temperature dependences:

cool to ~10 K @ ~5 K per minute

take one whole spectrum (in 1 s) every 0.5 K or 1 K.

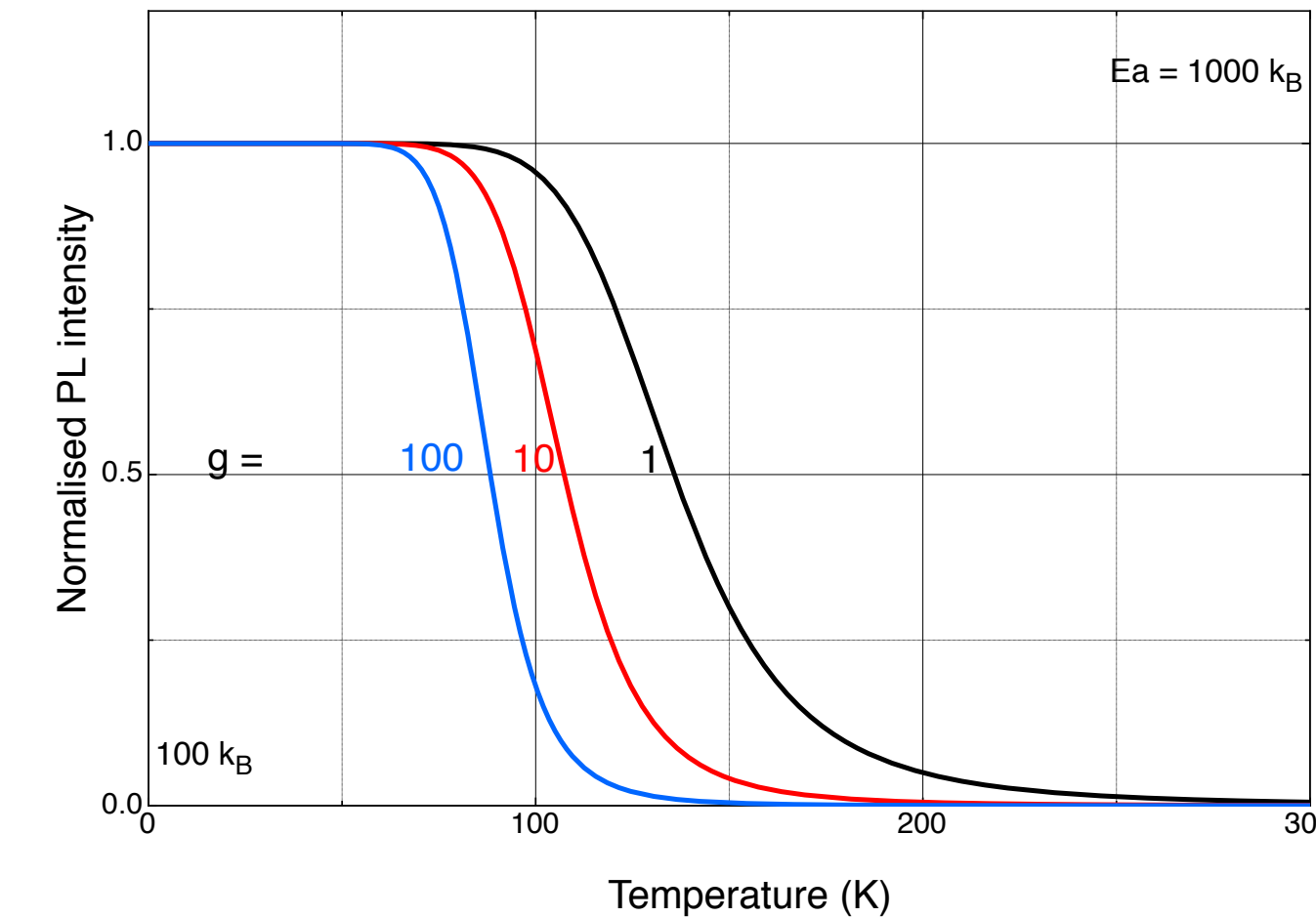
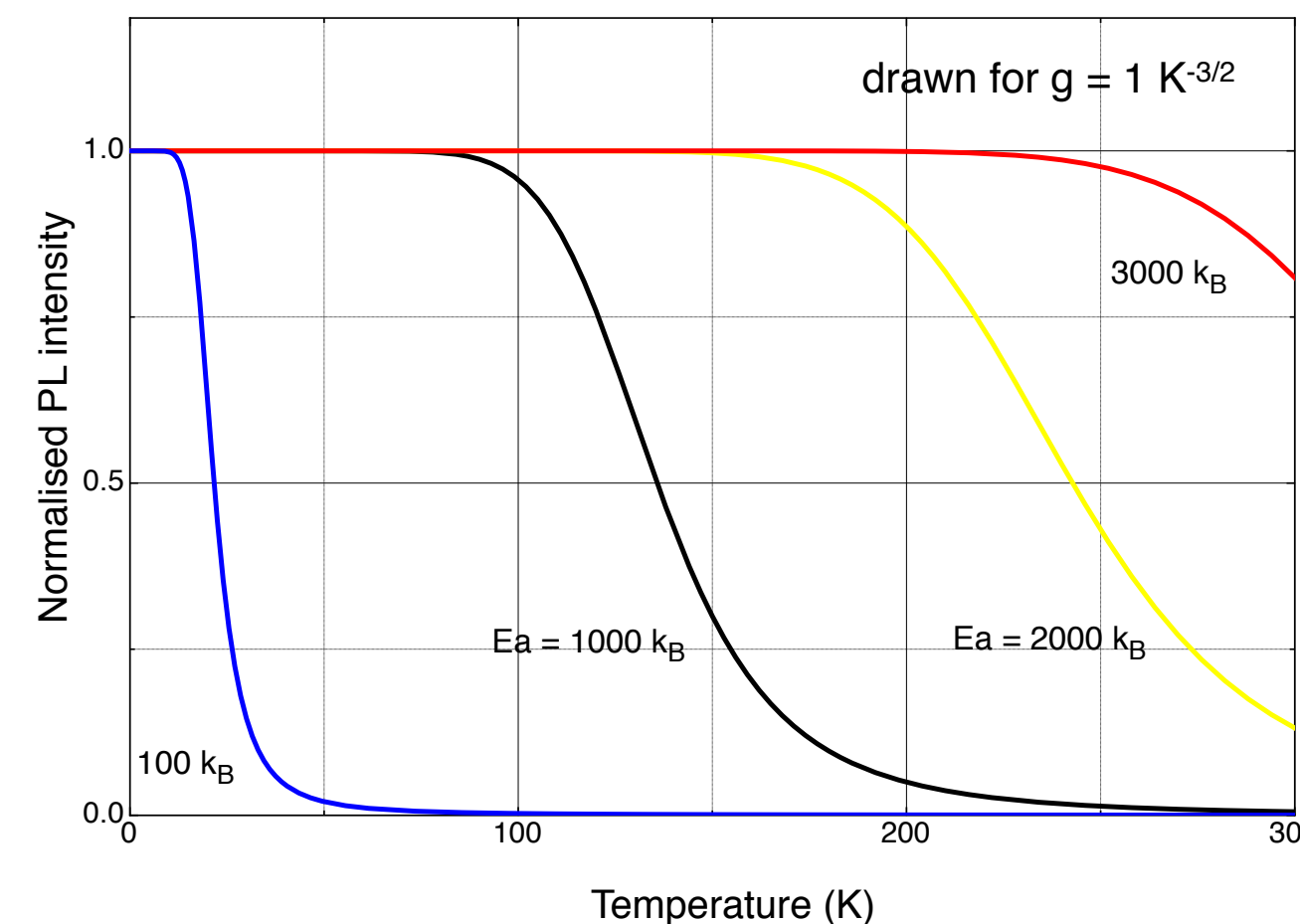
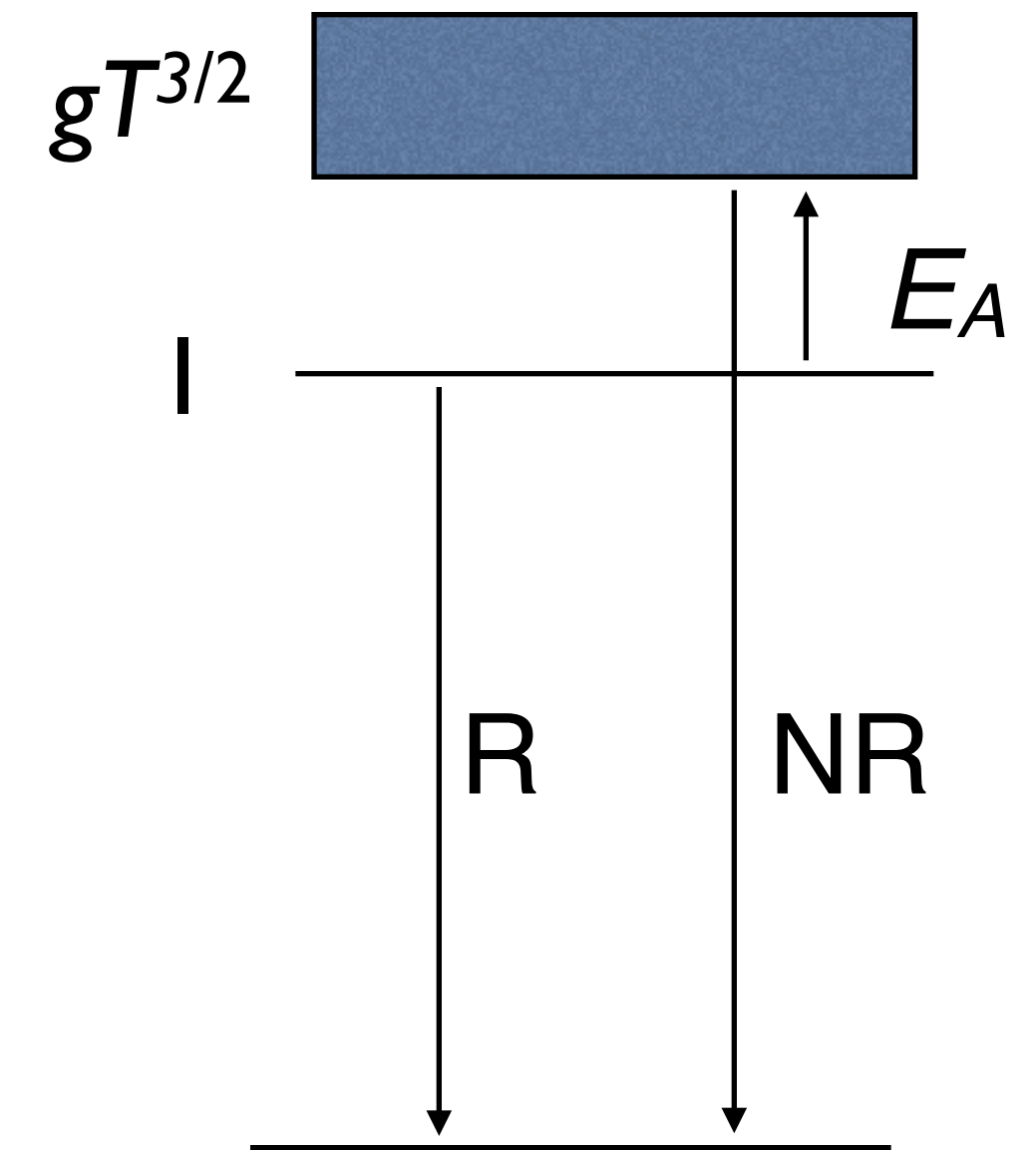
In the conventional model, an activation energy E_A separates a radiative (localised) state from a band of delocalised states, with a relative degeneracy ratio of $1:gT^{3/2}$.

The radiative fraction is then given by:

$$IQE(T) = \frac{I(T)}{I(0)} = \frac{1}{1 + gT^{3/2} \exp(-E_A / k_B T)}$$

which shows the well-known general form, with a nearly constant intensity up to a temperature $\sim E_A/10k_B$ followed by a fairly rapid decline as T increases further.

Note: $E_A = k_B T_A$; 1 eV \sim 11 600 K

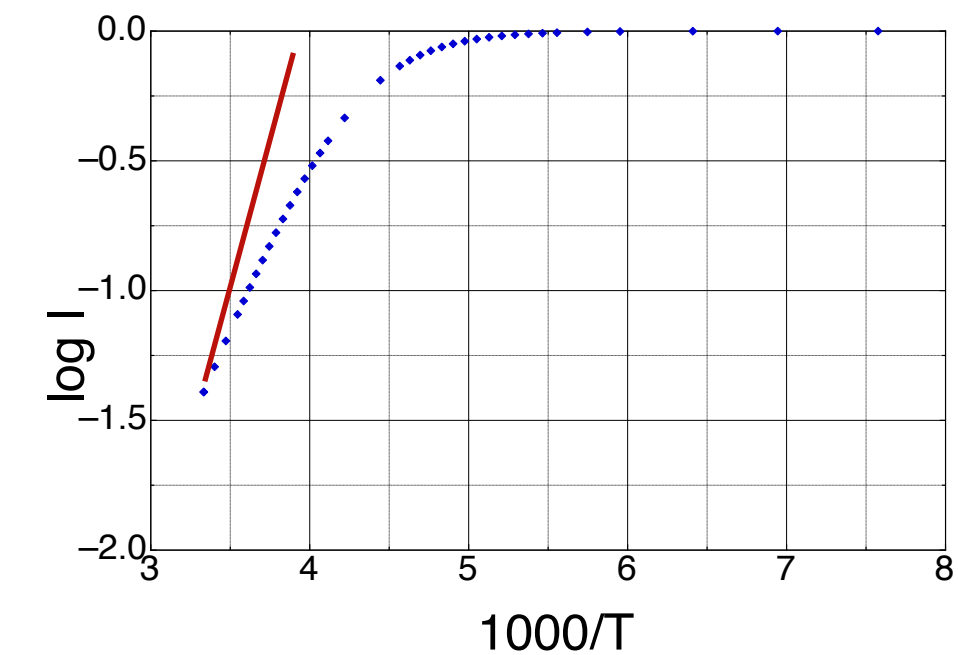


$$IQE(T) = \frac{I(T)}{I(0)} = \frac{1}{1 + gT^{3/2} \exp(-E_A / k_B T)}$$

- The $T^{3/2}$ dependence is relatively weak and is often ignored. It makes little practical difference to the goodness of fit for most data. Simplifying and approximating the formula* produces the Arrhenius form of the relation.:

$$\frac{I(T)}{I(0)} = \frac{1}{1 + G \exp(-E_A / k_B T)} \sim G^{-1} \exp(E_A / k_B T)$$

A plot of $\log(I)$ vs $1/T$ can be fitted *asymptotically* with a straight line to yield an *estimate* of E_A .



- $IQE(T = 300 \text{ K})$ is the value most often quoted, but the formula carries the implication that all semiconductor luminescence is 100% efficient at low temperatures: this is unlikely to be true.

- We can usefully define the *half-power point* at which temperature the luminescence intensity drops to half-maximum:

$$T_{1/2} = \frac{E_A}{k_B \ln G}$$

This parameter provides a ‘sanity check’ on the fitness of the model (see later).

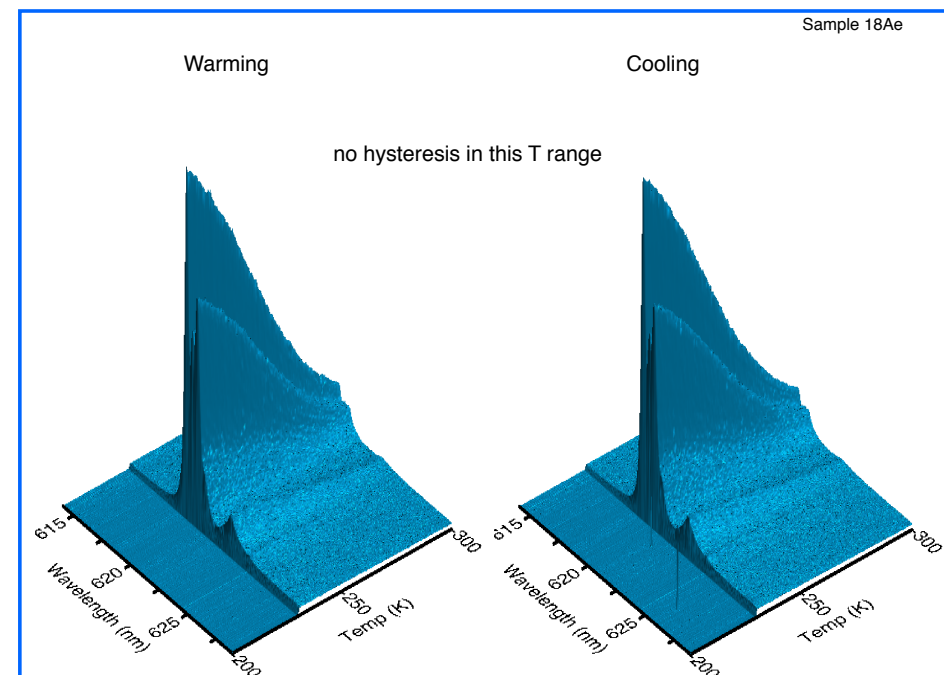
*as is my wont

GaN:Mg template

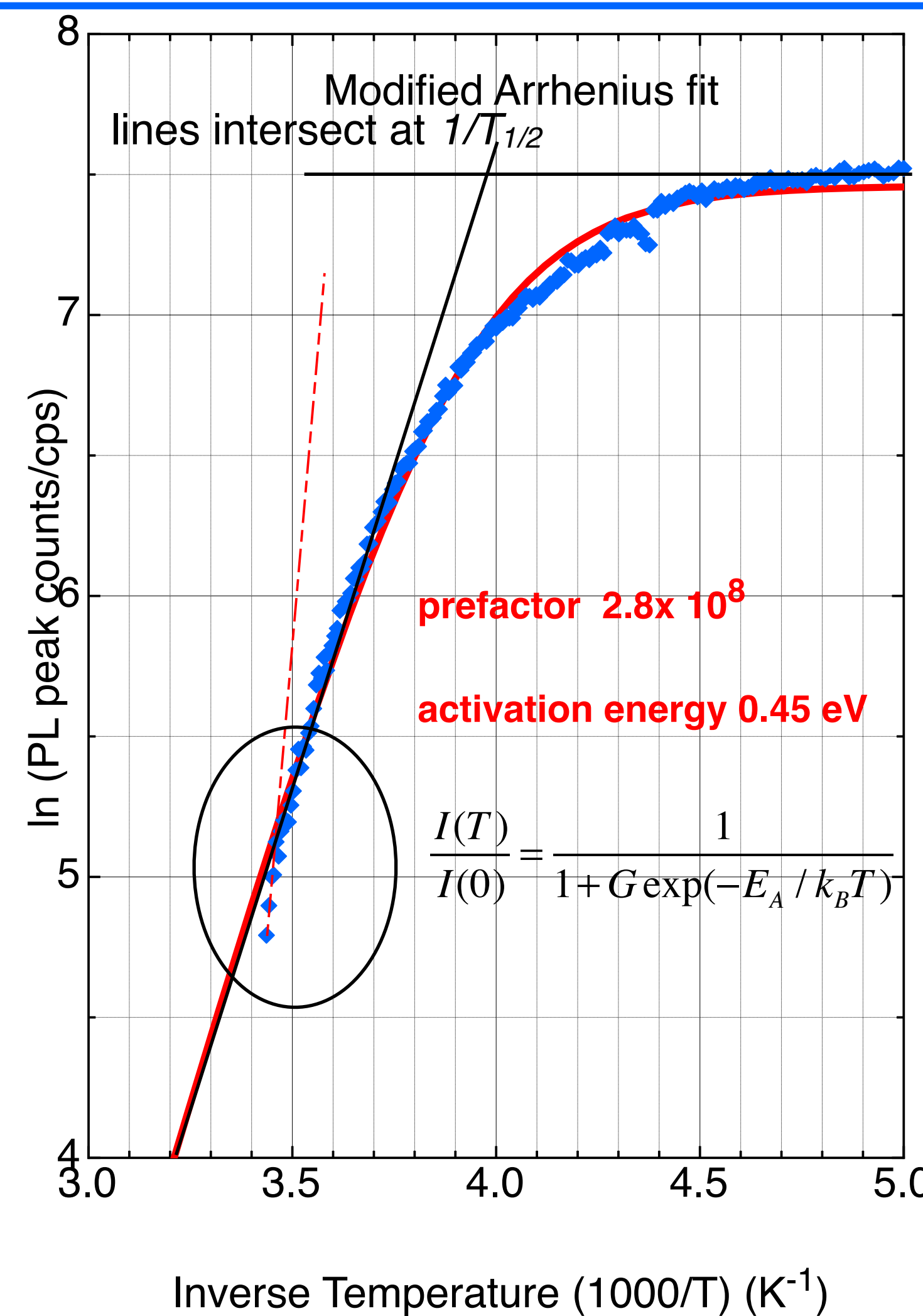
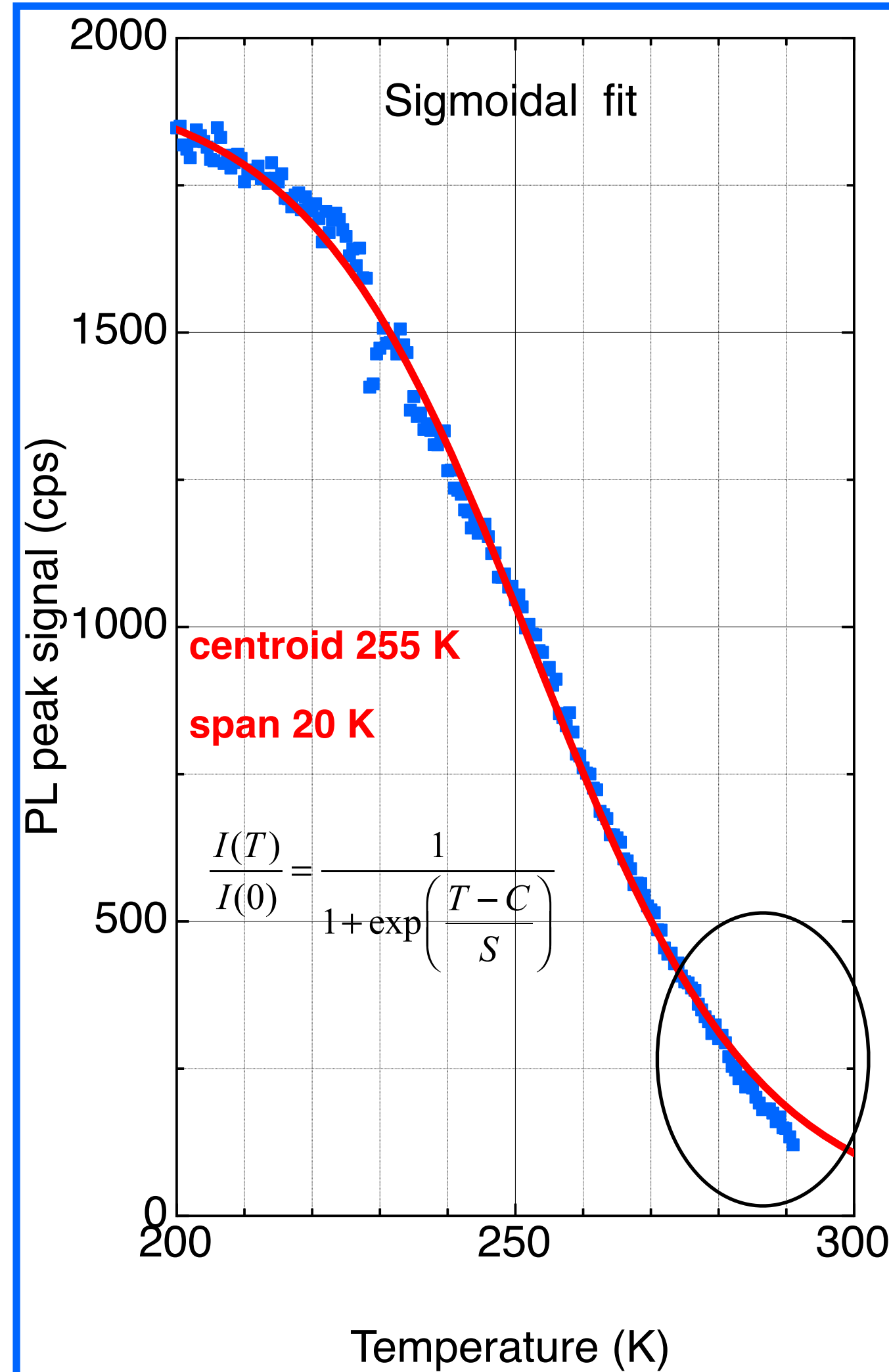
Eu ion implantation

HTHP annealing

Defect Eu0 is dominant near room temperature.



5D_0 to 7F_2 transition



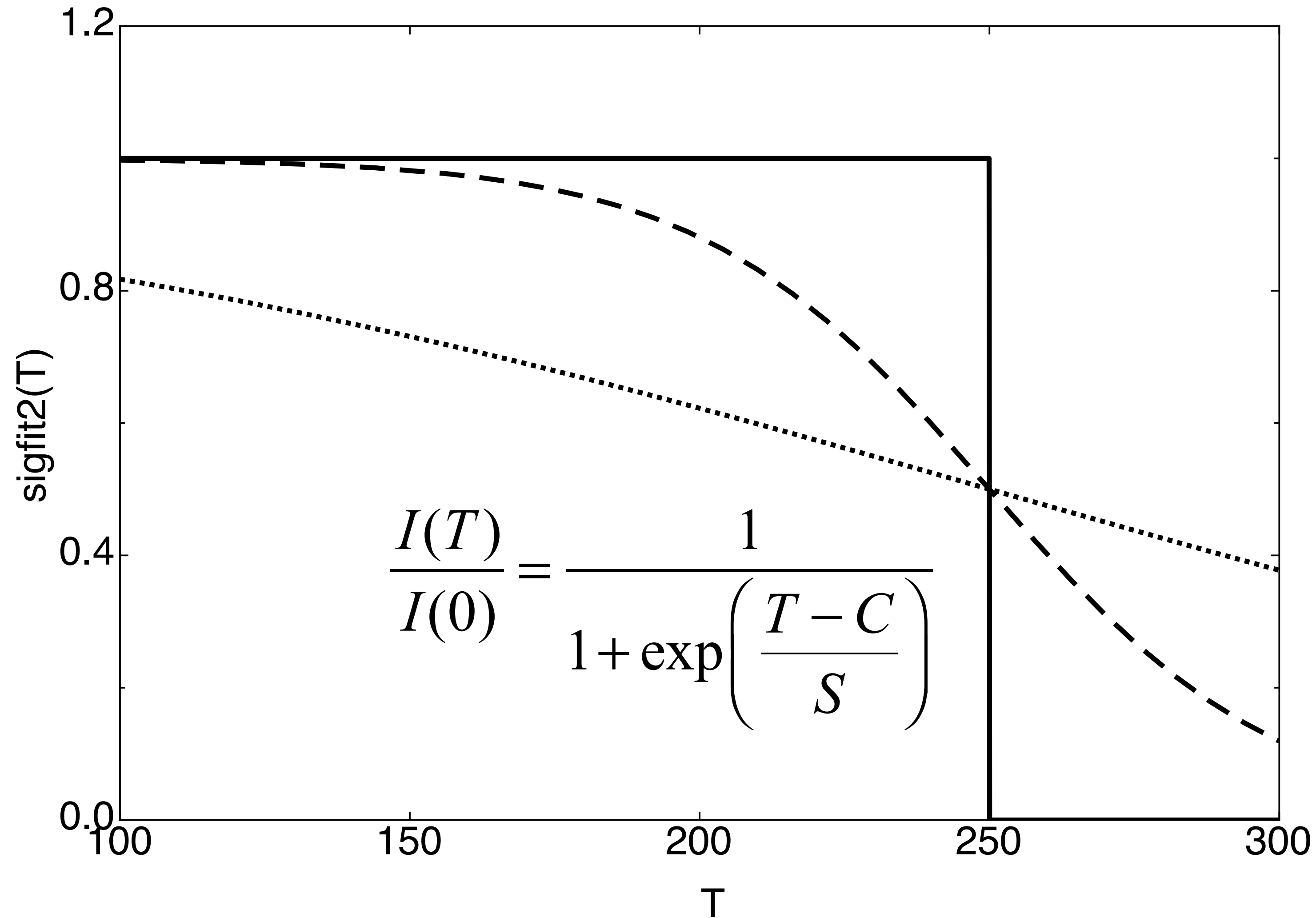
K. P. O'Donnell, P. R. Edwards,
M. J. Kappers, K. Lorenz, E. Alves
and M. Boćkowski

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All fits tend to deviate at low PL intensity, due partly to errors in background subtraction.



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The temperature dependence of photoluminescence in a-Si: H alloys

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Abstract

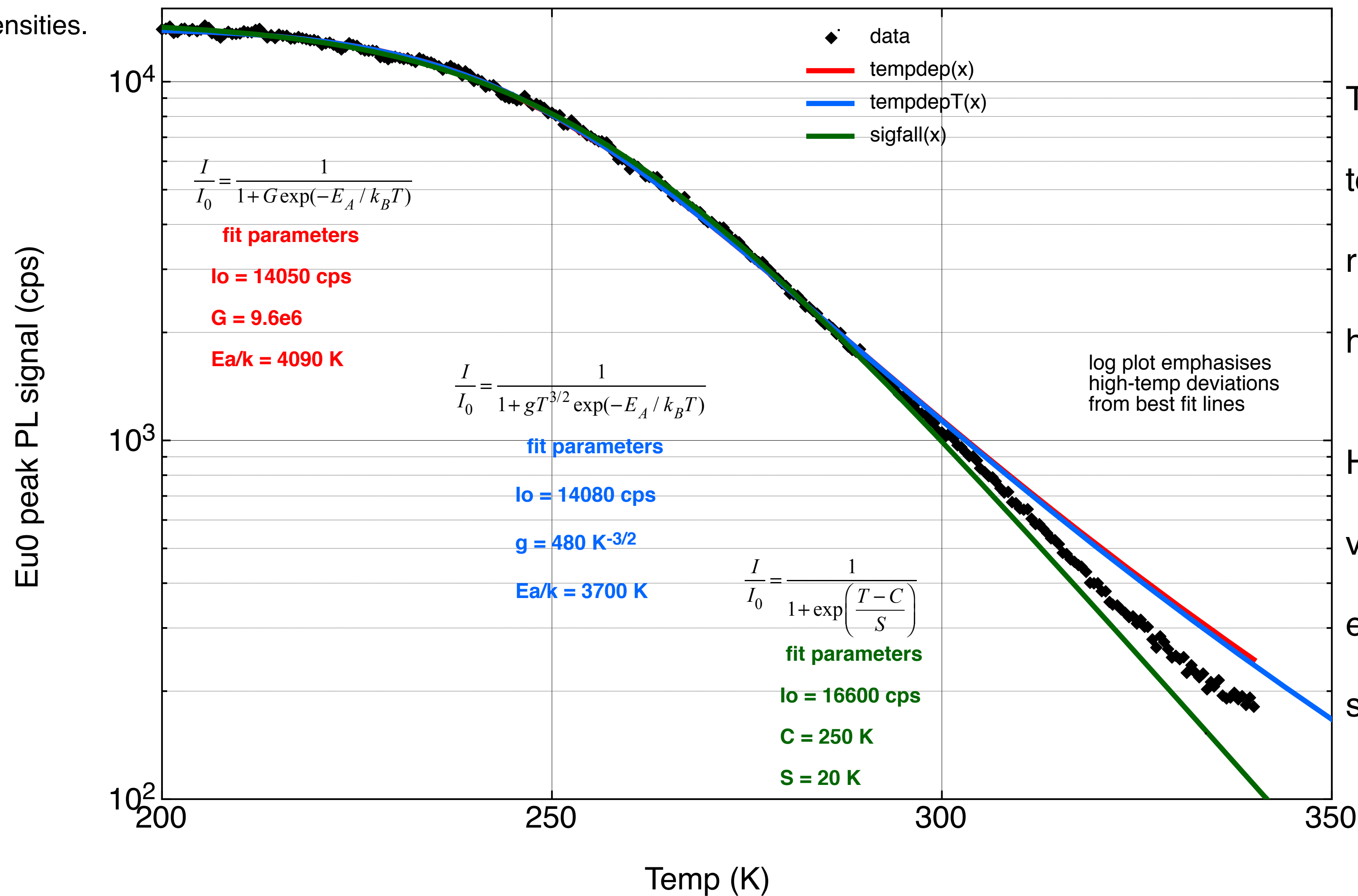
Photoluminescence intensity observed near 1.3 eV in sputtered a-Si : H has been measured as a function of temperature for several samples prepared under differing conditions. The data are shown to obey an expression derived from a law of the form

$$\frac{\rho_{NR}}{\rho_R} \sim e^{T/T_0}$$

where p_{nr} and p_r are the probabilities for non-radiative and radiative recombination. We find $T_0 \approx 23$ K independent of sample preparation conditions.

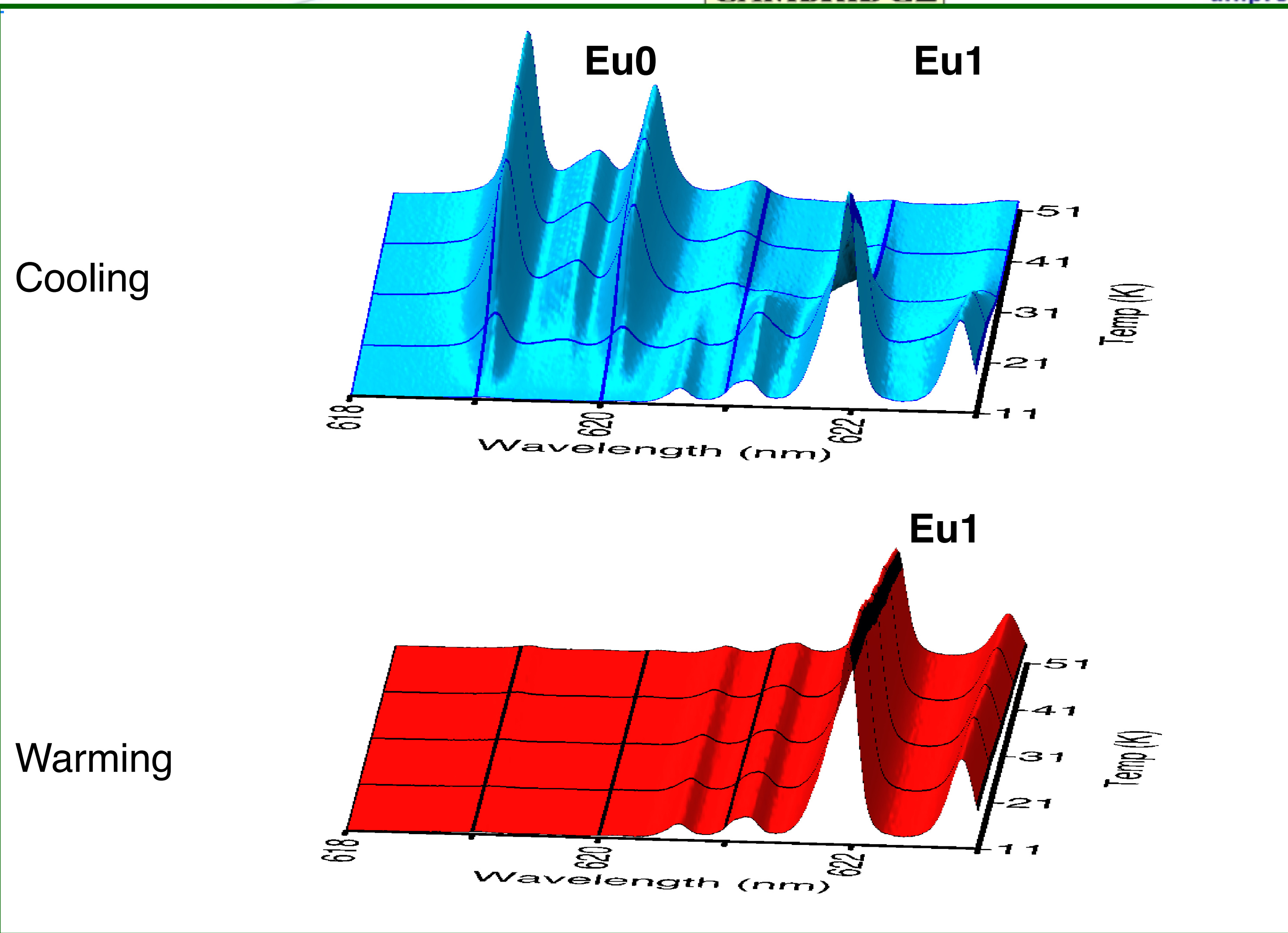
The temp dependences are similar since the radiative fraction $\frac{\rho_R}{\rho_R + \rho_{NR}} \sim \frac{1}{1 + \exp\left(\frac{T}{T_0}\right)}$ is a sigmoid with $T_{1/2} = 0$ K and $S = T_0$

Note restricted range of intensities.



The fits are naturally weighted towards the low temperature region where intensities are higher.

High temperature intensity values suffer more from errors in background subtraction.

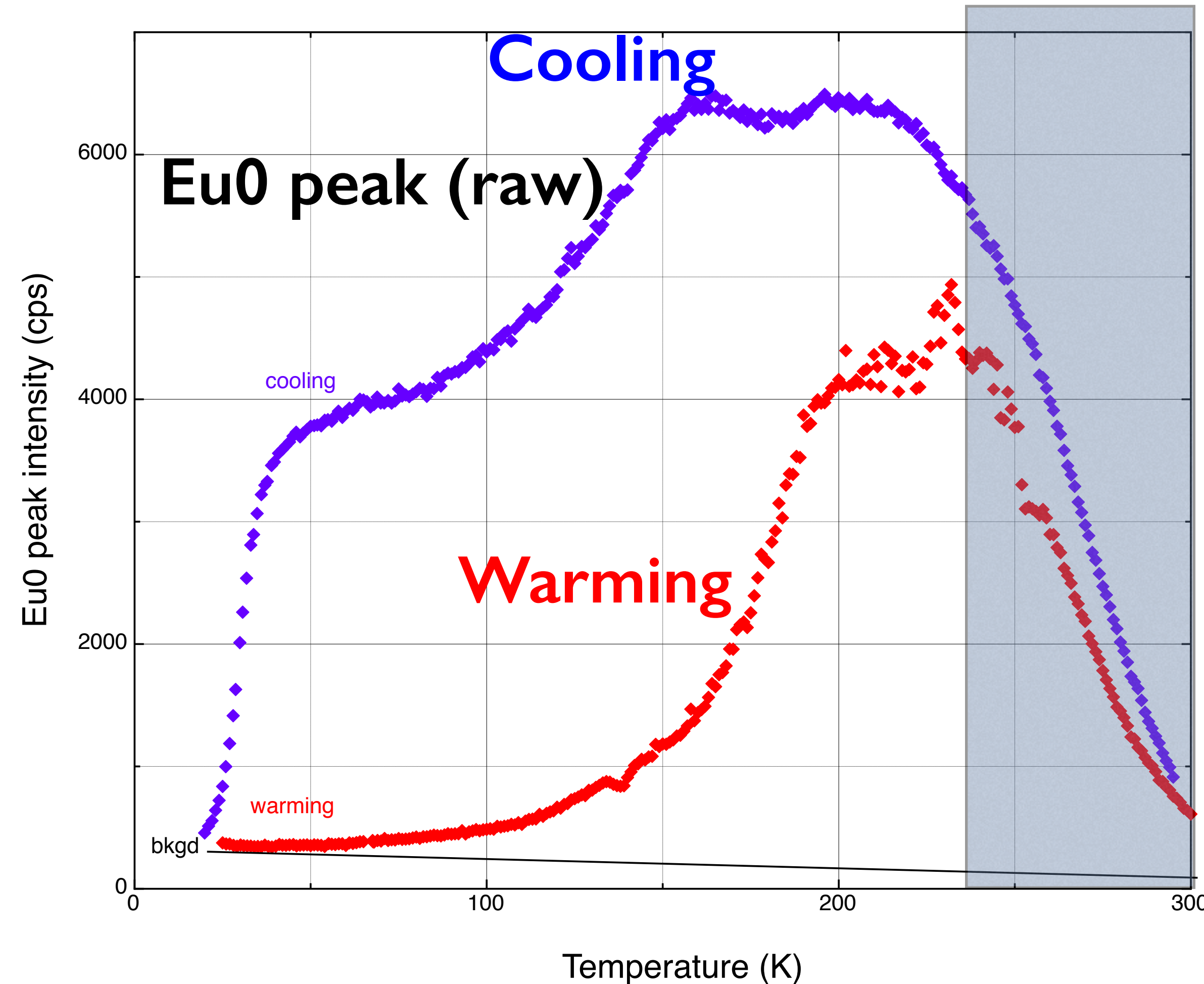


Eu0 anomalous;
Eu1 quasi-normal.

Eu0 absent;
Eu1 normal.

Sample 3a complete switching

The anomalous (hysteretic) behaviour here can be described by the usual equations if we allow negative energy as a fit parameter.



The shaded region has a normal temp. dependence: whether warming or cooling, the intensity is lower at higher temperature.

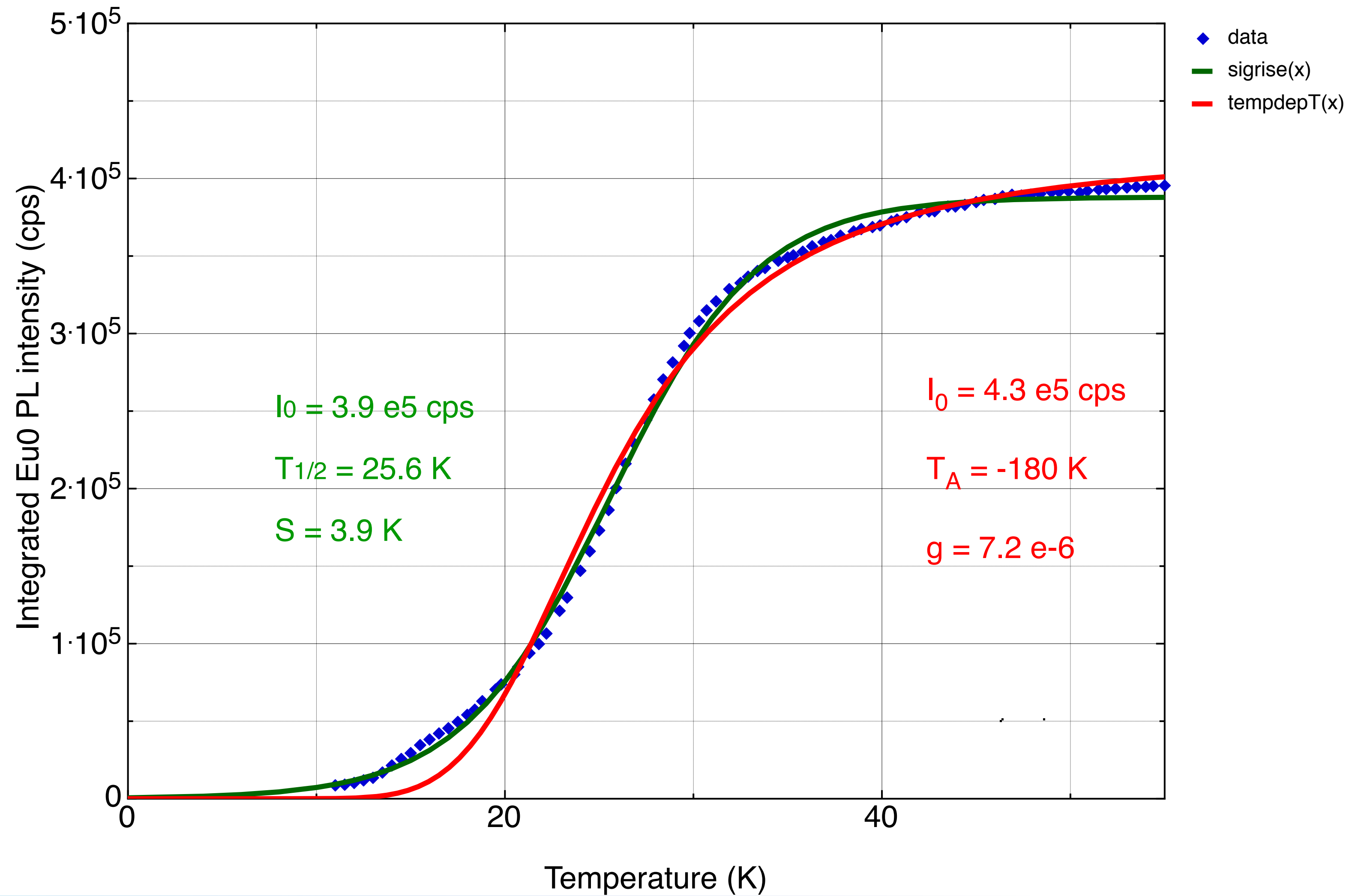
A rising step is just the complement of a falling step- Sigmoid Frond

$$1 - \frac{1}{1 + gT^{3/2} \exp\left(\frac{-E_A}{k_B T}\right)} = \frac{1}{1 + (1 / gT^{3/2}) \exp\left(\frac{E_A}{k_B T}\right)}$$

$$1 - \frac{1}{1 + \exp\left(\frac{T_{1/2} - T}{S}\right)} = \frac{1}{1 + \exp\left(-\left(\frac{T_{1/2} - T}{S}\right)\right)}$$

fall

rise



Odd behaviour of the BL band in high-resistivity GaN:Zn: **power-tuneable** thermal quenching

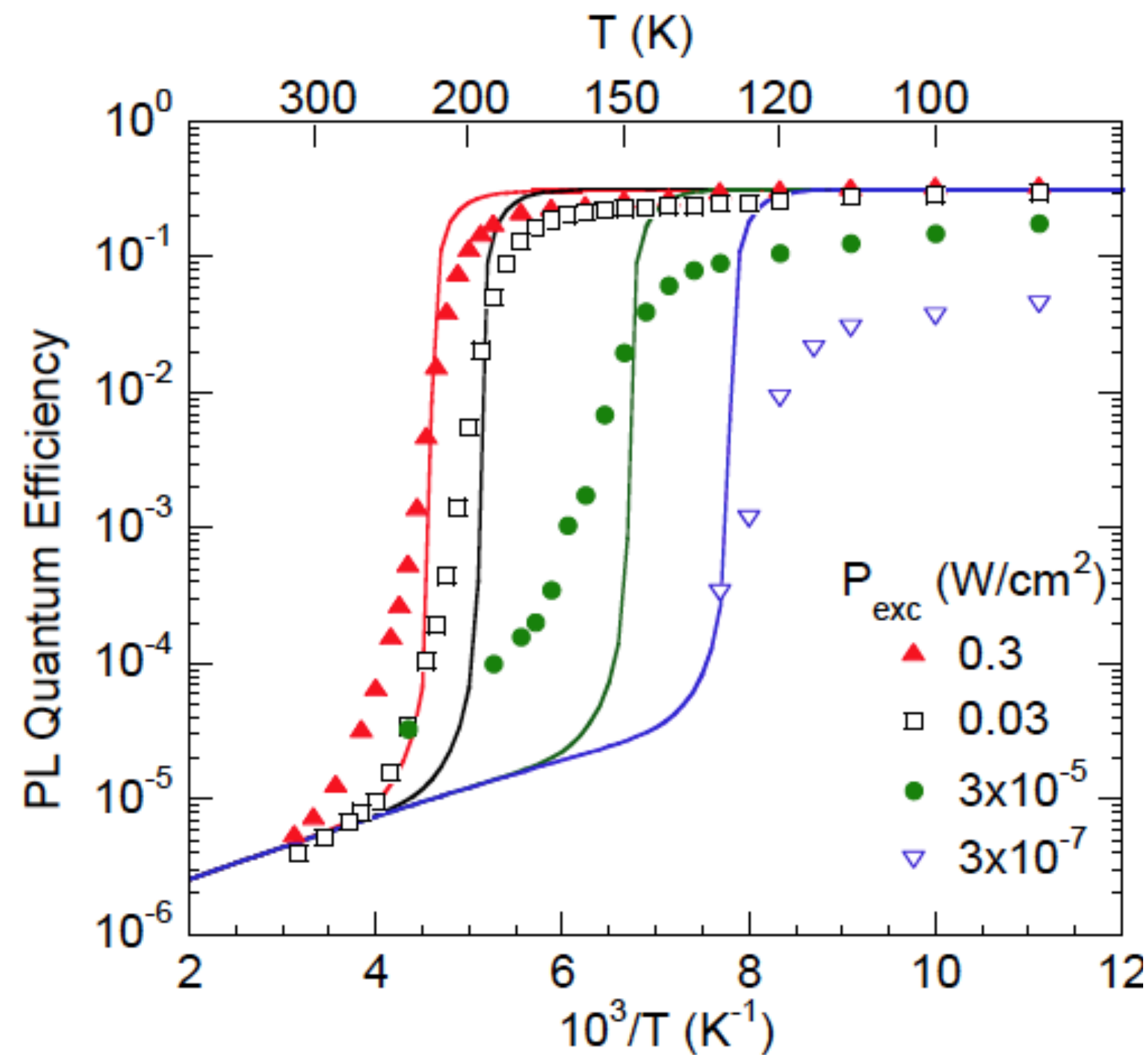


Figure 2 Temperature dependence of the quantum efficiency of the BL band in high-resistivity Zn-doped GaN for selected P_{exc} .

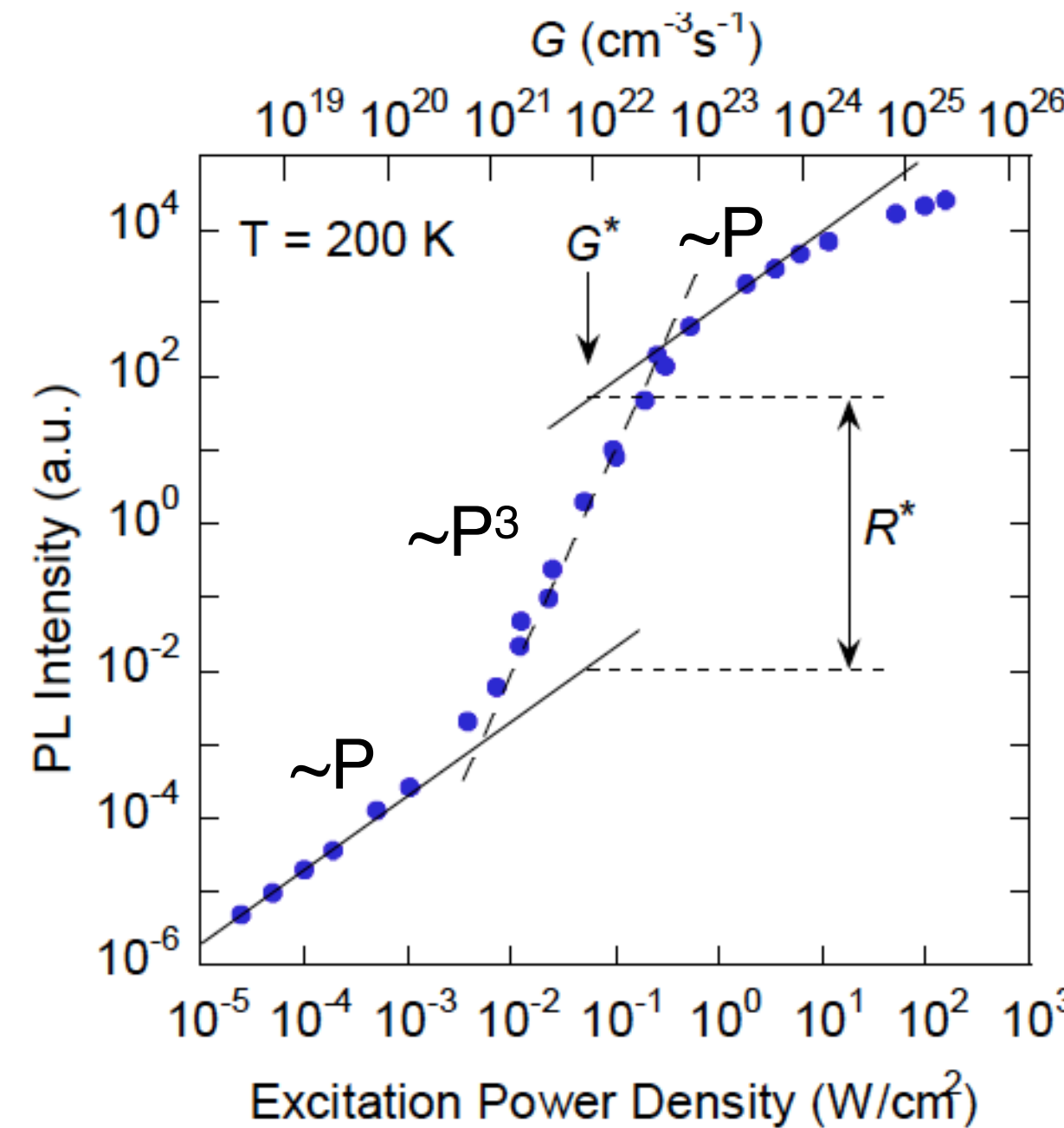
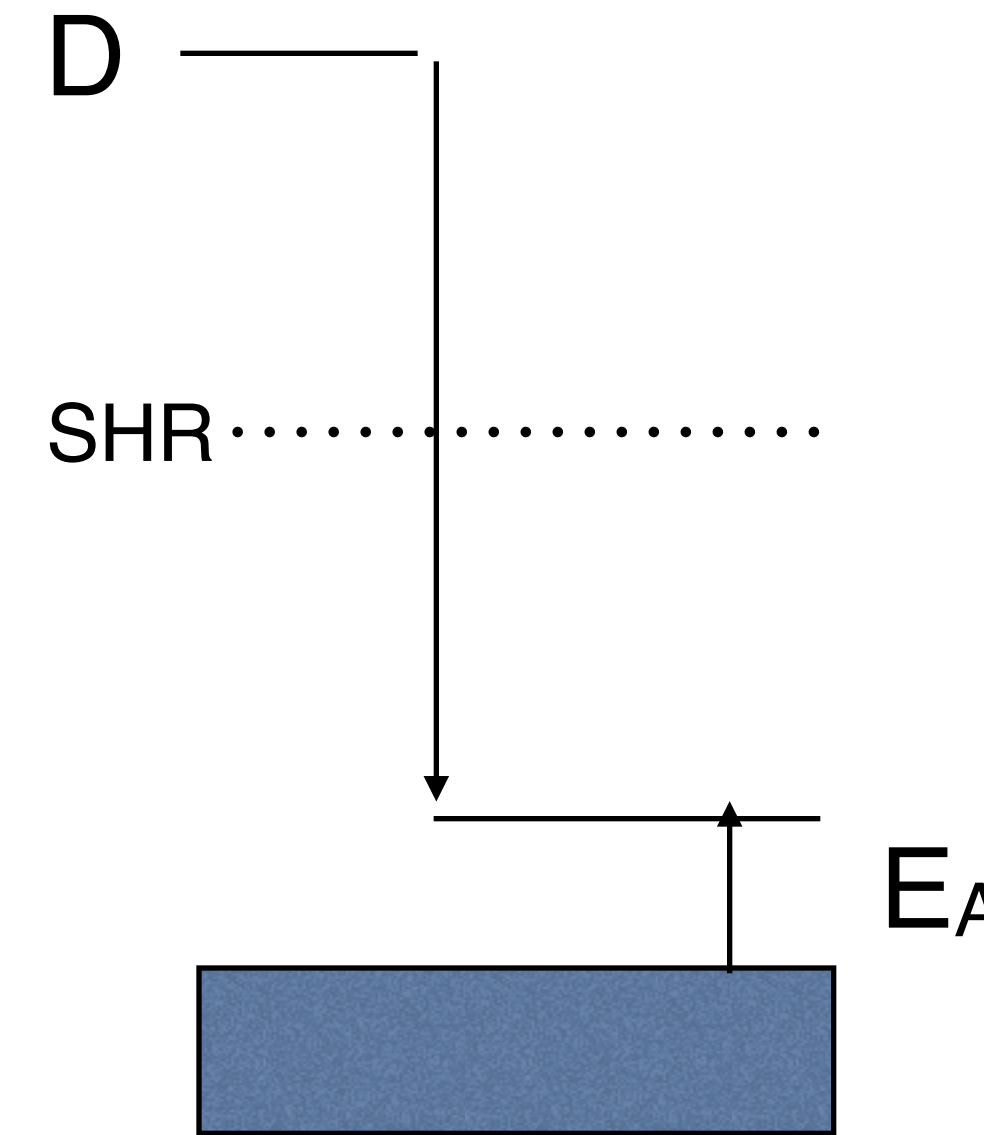
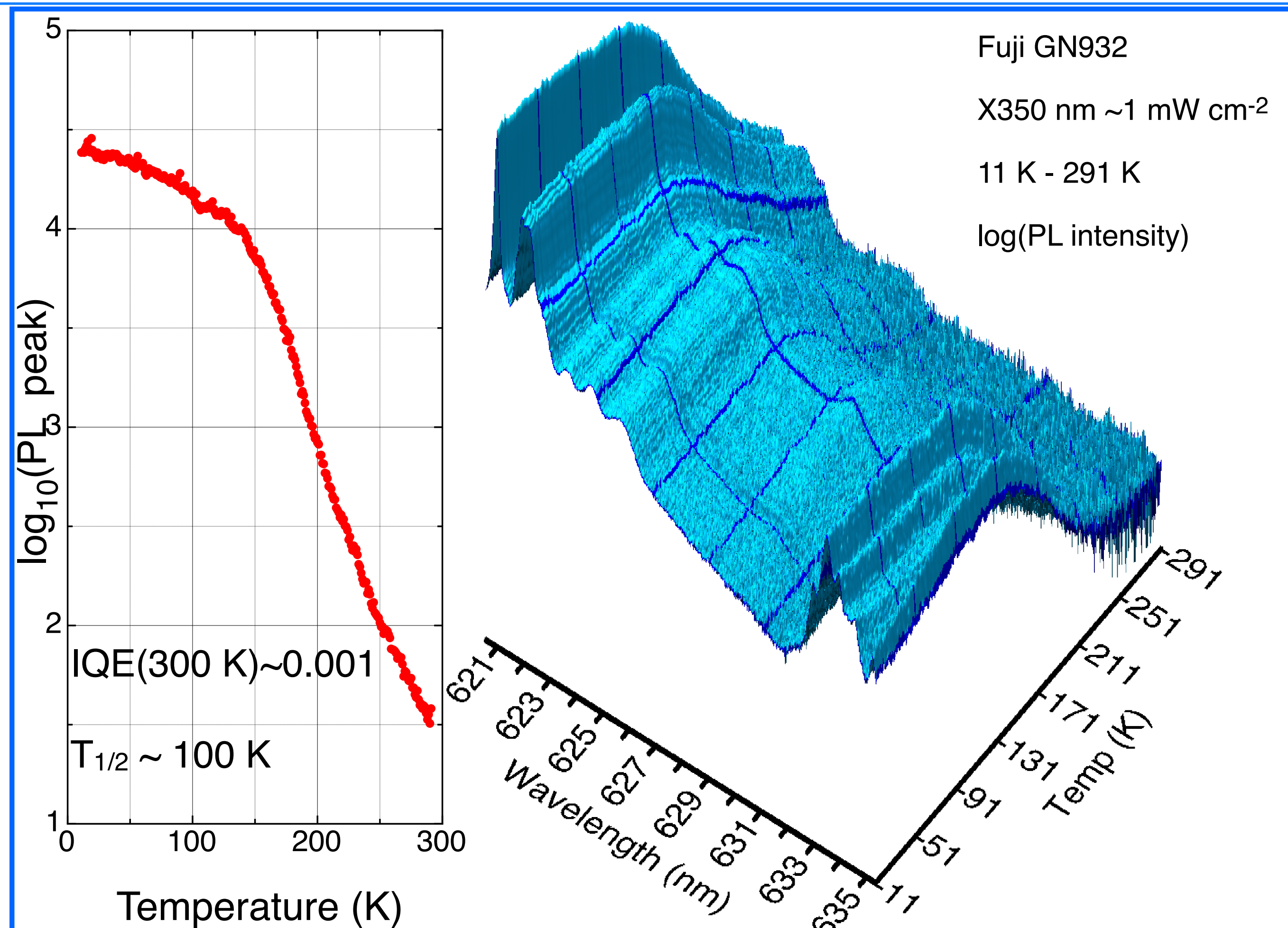


Figure 4 Dependence of the BL band intensity on the excitation intensity at $T = 200$ K.

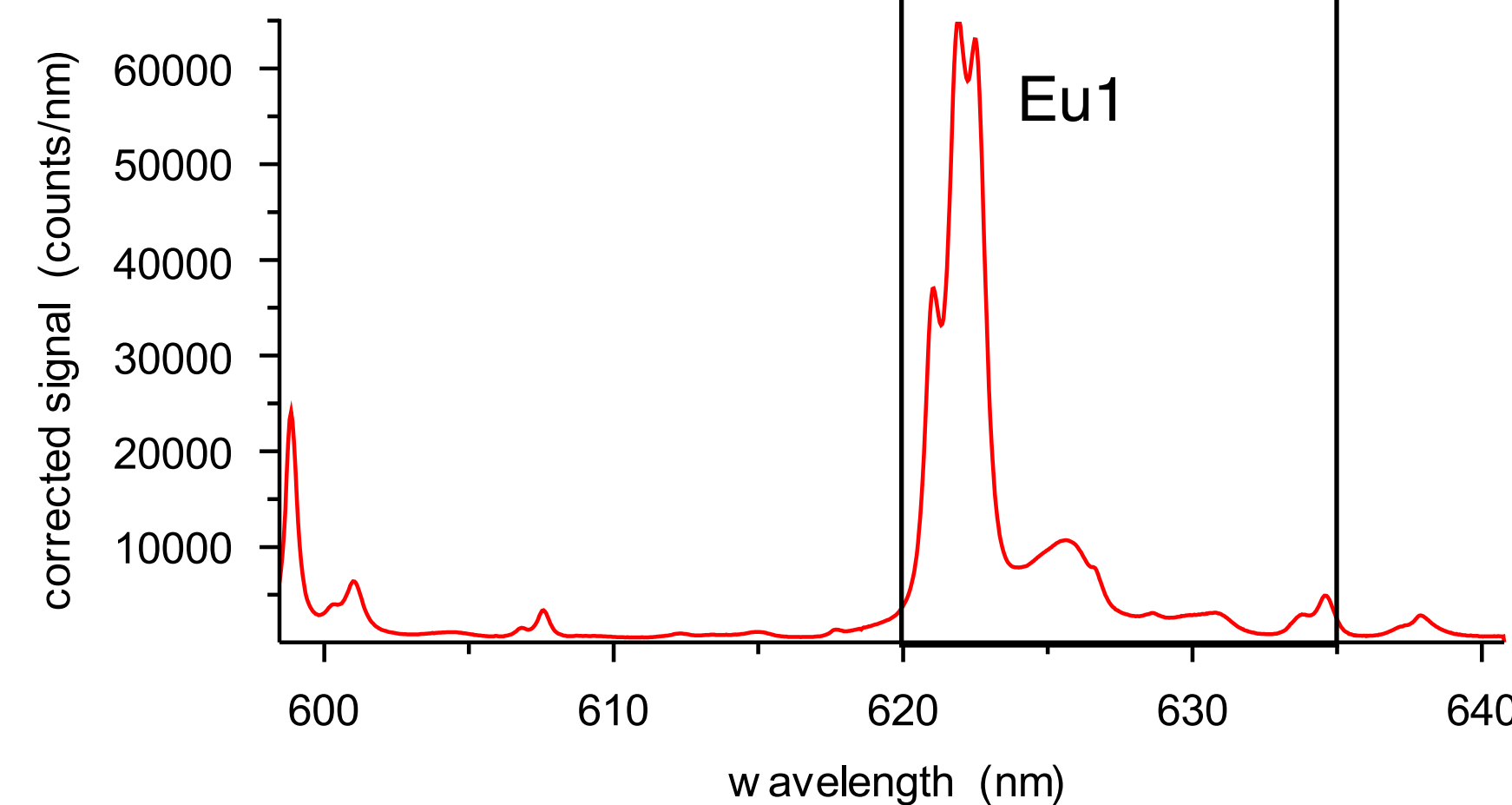
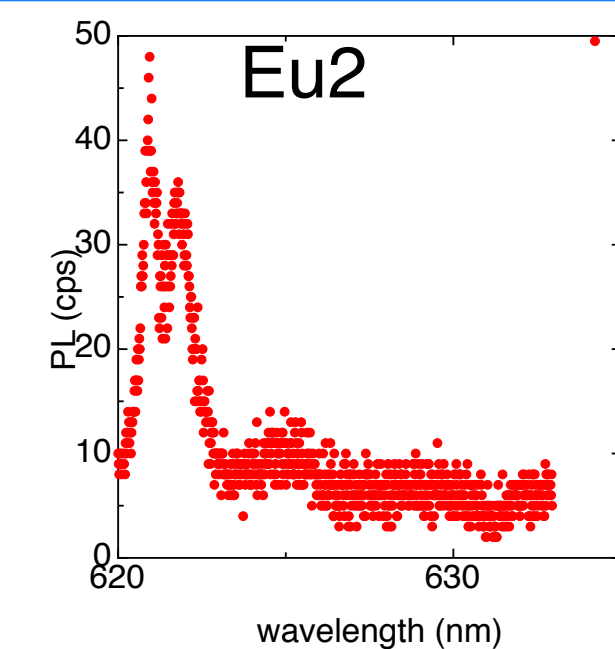


GaN:Zn (deep acceptor) with shallow donor and *Shockley-Hall-Reed* e-h recombination centre.

Phys. Status Solidi C 10, No. 3, 515–518 (2013)

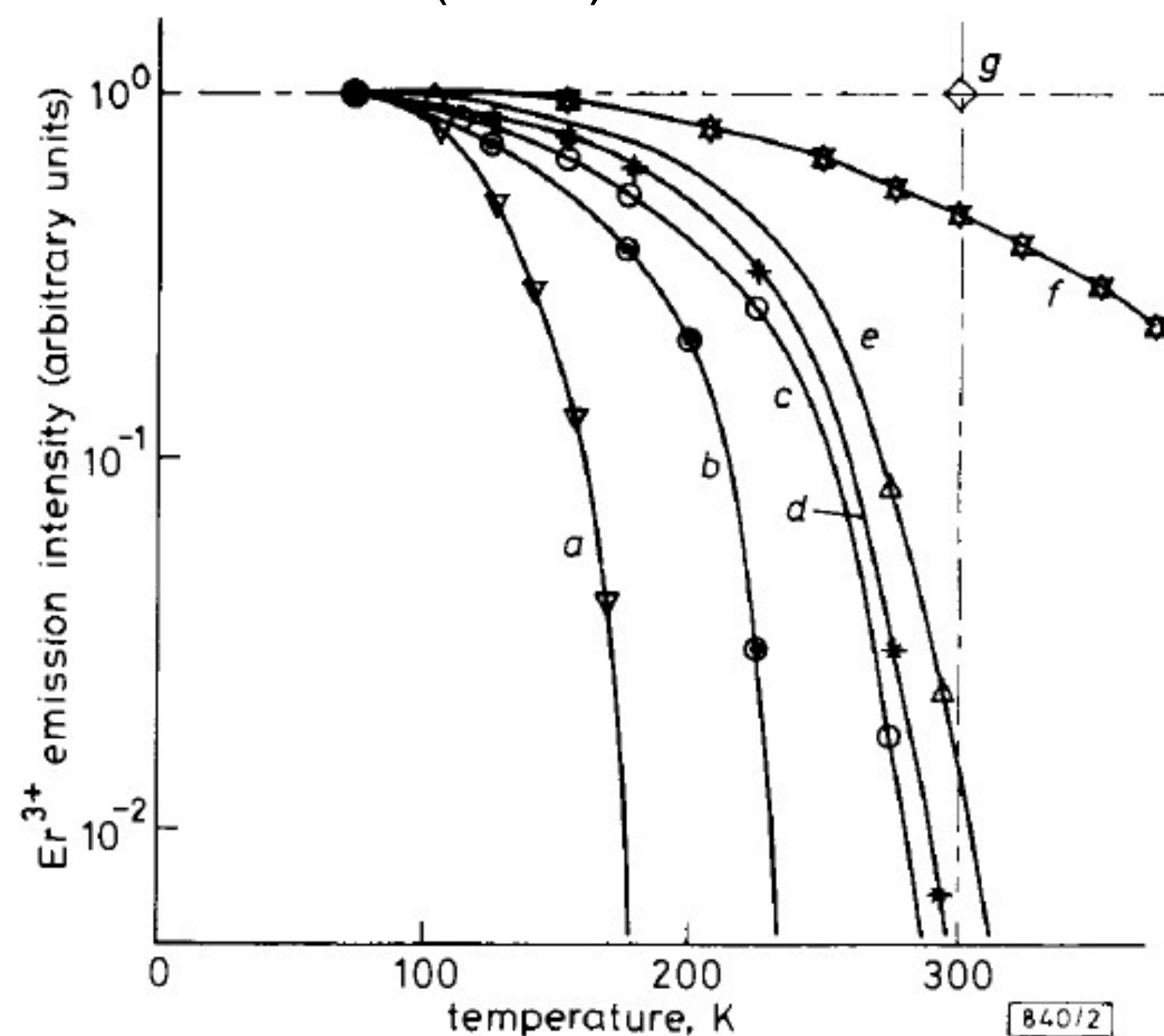


Low power photoexcitation-
room temp. spectrum



High power cathodoexcitation- room temp. spectrum

P.N. Favennec et al. (1989)


Fig. 2 Er^{3+} emission intensity against host semiconductor temperature

 Materials are implanted with Er ions: $E = 330 \text{ keV}$, $\phi = 10^{13} \text{ Er}^+$ cm^{-2} . E_G values are given at room temperature

 a $\text{Ga}_{0.38}\text{In}_{0.62}\text{As}_{0.84}\text{P}_{0.12}$ ($E_G = 0.807 \text{ eV}$)

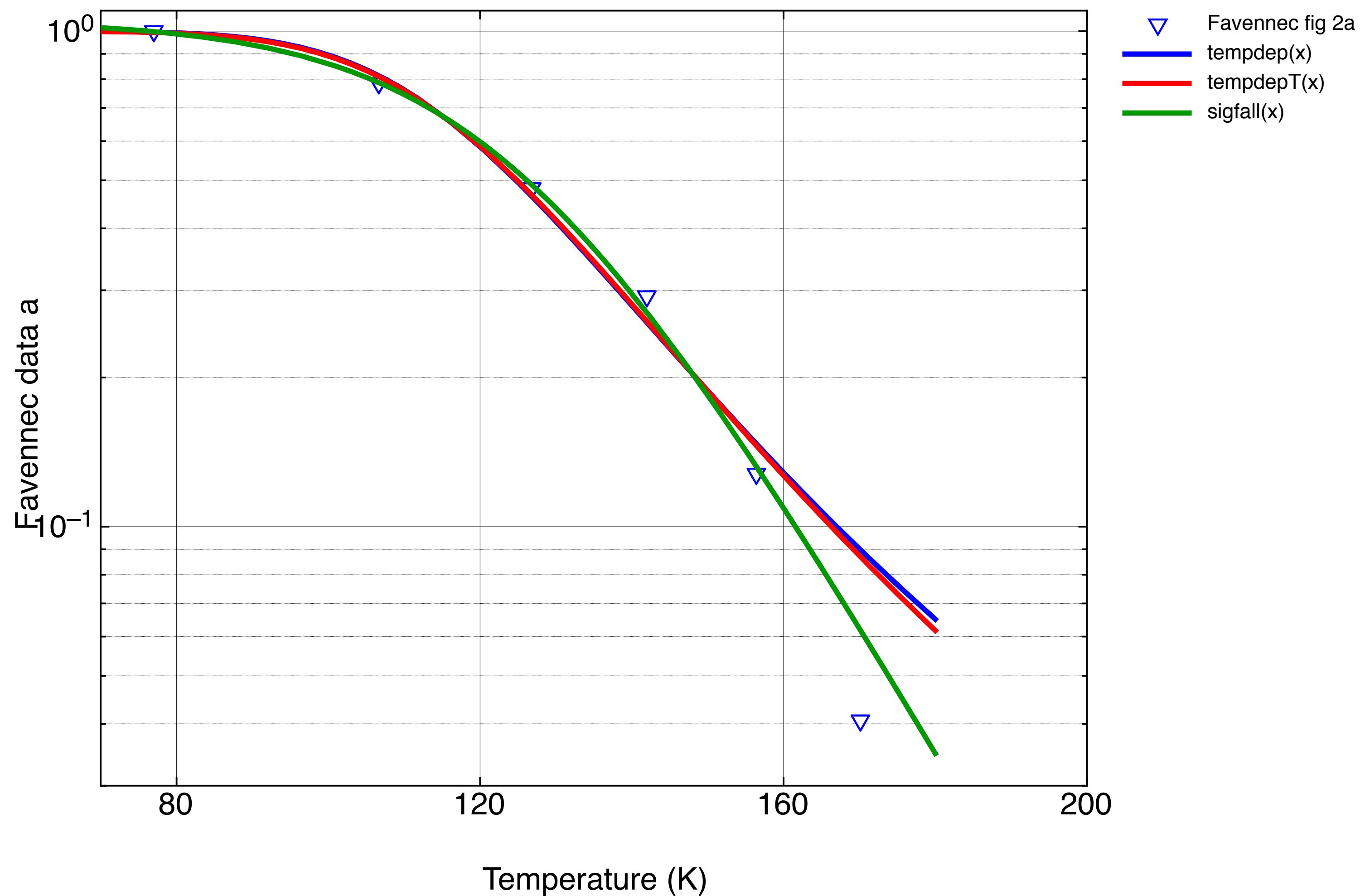
 b Si ($E_G = 1.12 \text{ eV}$)

 c InP ($E_G = 1.27 \text{ eV}$)

 d GaAs ($E_G = 1.43 \text{ eV}$)

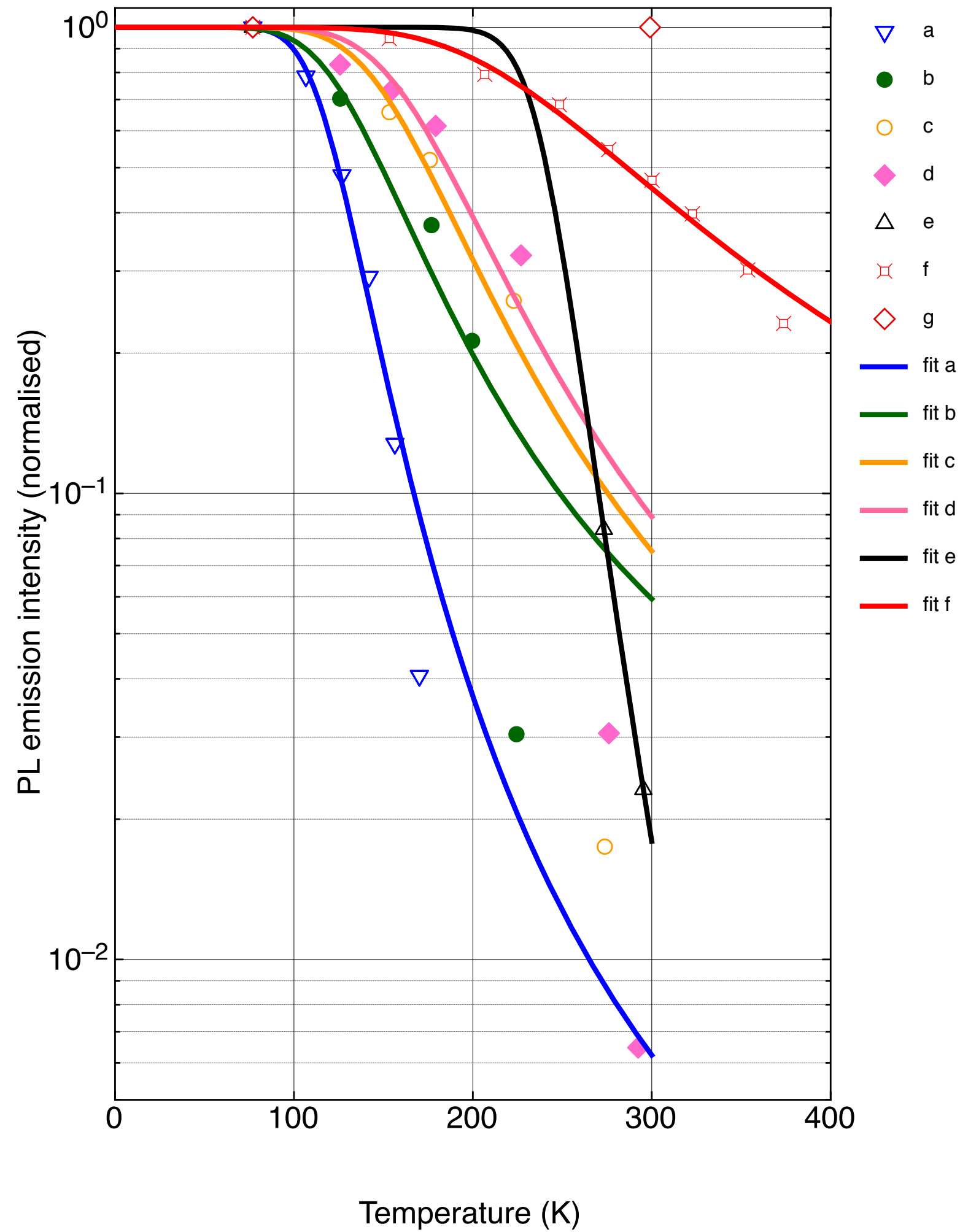
 e $\text{Al}_{0.17}\text{Ga}_{0.83}\text{As}$ ($E_G = 1.67 \text{ eV}$)

 f ZnTe ($E_G = 2.26 \text{ eV}$)

 g CdS ($E_G = 2.42 \text{ eV}$)


ELECTRONICS LETTERS 25th May 1989 Vol. 25 No. 11

After Favennec et al (1989)

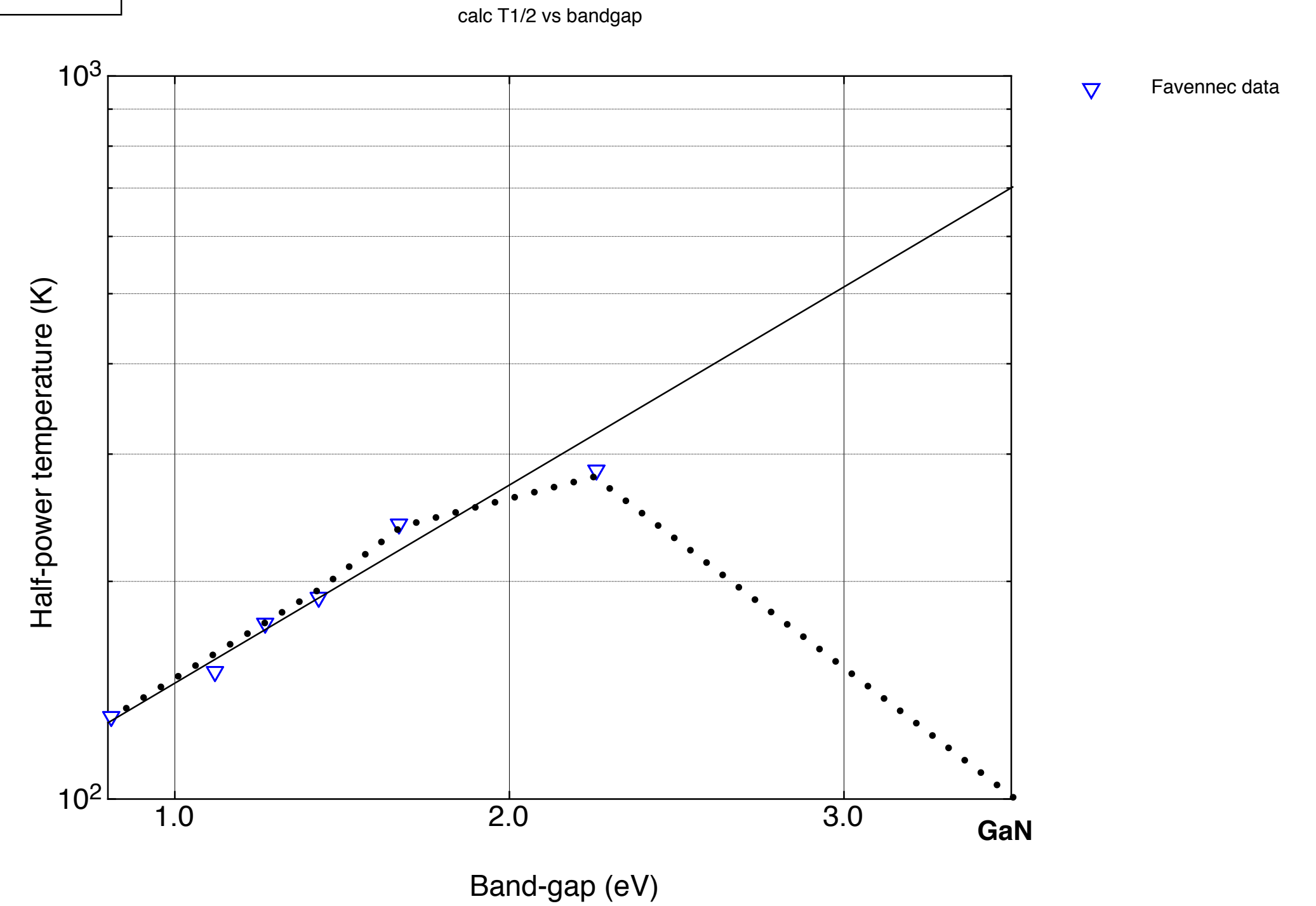


sample	semic. gap	G	Ta (K)	T1/2 (K)
a	GaInAsP 0.807 eV	5900	1100	130
b	Si 1.12 eV	240	820	150
c	InP 1.27 eV	400	1050	175
d	GaAs 1.43 eV	440	1130	190
e	AlGaAs 1.67 eV	7.50E+08	4900	240
f	ZnTe 2.26 eV	65	1190	285
g	CdS 2.42 eV	-	-	

LUMINESCENCE OF ERBIUM IMPLANTED
IN VARIOUS SEMICONDUCTORS: IV, III-V
AND II-VI MATERIALS

11th April 1989

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- The temperature dependence of luminescence from semiconductors is well described by the conventional R-NR model, unless it isn't.
- A parameter that can be extracted from the data is E_A , the activation/localisation/binding energy (delete as appropriate).
- More useful in a practical sense is the IQE (RT), but you may prefer to quote the half-power temperature $T_{1/2}$.
- In terms of the fitting parameters for the modified Arrhenius fit, $T_{1/2} = \frac{E_A}{k_B \ln G}$; for the sigmoidal fit $T_{1/2} = T_{1/2}$. (The derivation of an expression for $T_{1/2}$ in the conventional R-NR model is left as an exercise for ~~Phil Dawson~~ students.)
- In Reschikov's work, a slight modification of the simple temperature-dependence model produces wonderful complications.
- A reanalysis of Favennec's data 25 years on teaches us that extrapolation is dangerous.